

# THE MAGNETOSPHERIC CUSP: STRUCTURE AND DYNAMICS

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## ABSTRACT

Understanding the polar cusps is essential for a thorough understanding of the entire physics of the magnetosphere, and of the dynamical interaction between the solar wind and any planetary magnetosphere. Energetic electrons are unique to fully assess magnetic field-line topology and thus should be able to clearly delineate regions of open and closed magnetic field lines in the high latitude regions and contributed crucially to understanding and resolving an internal debate going on between groups measuring only the lower energy ( $< 20$  keV) plasma. Energetic electrons with high and stable flux were observed in the high latitude boundary/cusp region when the IMF had a predominate positive  $B_z$  component. With measurements at larger separations and more coordination of multiple satellite measurements for particular cusp crossings will it become more evident what the true nature of the cusp is and what roles the cusps play. The boundary normal, velocity and timing analysis obtained by all four spacecraft indicates that the multiple cusp phenomena is most likely caused by the oscillation of the single northern cusp which was shifted back and forth. Cusp oscillating with a period of 22 min are observed by Cluster satellite in the high latitude region, in the meantime, the cold-dense plasma with fluctuations (20 min period) are observed in the dusk-side of the tail plasma sheet by Geotail satellite. This is consistent with the idea that the high latitude reconnection during northward IMF is the responsible mechanism of the formation of the cold-dense plasma sheet.

## 1. BRIEF HISTORY ON THE CUSP

- *~Qin dynasty (221-206 B.C.) Magnetic compass discovered in China. The first person recorded to have used the compass as a navigational aid was Zheng He (1371-1435),*

*from the Yunnan province in China, who made seven ocean voyages between 1405 and 1433.*

- *1600 William Gilbert publishes in London "De Magnete" ("on the magnet"). His explanation of the compass: the Earth is a giant magnet.*
- *Maxwell (~1880) showed that a perfect conductor adjacent to a dipole formed an image dipole*
- *Chapman and Ferraro (1931) first induced the basic nature of the Earth's magnetosphere, its 2-D and 3-D topology has indicated the existence of a dayside magnetic cusp.*
- *Spreiter and Summers (1962) predicted a stagnation flow in the cusp region by using a gas dynamics model*
- *Heikkila and Winningham (1971) and Frank (1971) showed a high-latitude band of low-energy particle precipitation with magnetosheath-like properties on the dayside at low altitudes which have been accepted as the first evidence to discover the magnetospheric cusp.*

## 2. INTRODUCTION TO THE CUSP

The boundaries of the magnetosphere including the polar cusp are key regions for the transfer of mass, momentum and energy from the solar wind into the magnetosphere no matter IMF is southward or northward. The first identification of a thin layer of magnetosheath plasma located immediately inside the magnetopause was made by Hones [1972] who also

introduced the term "Boundary Layer". Since then, the morphological characteristics as well as plasma properties of the magnetospheric boundary layer have been studied rather intensively [Rosenbauer et al., 1975; Eastman et al., 1976; Haerendel, 1978; Lundin et al., 1985a, 85b; Newell et al., 1988, 1998]. The term "Low-latitude boundary layer", was apparently introduced by Haerendel [1978] to distinguish the very different properties observed at latitudes below about  $50^\circ - 60^\circ$  on the magnetopause surface. In addition there are three more boundary regions (in the high latitude) which are assumed to connect directly to the magnetosheath: They are the plasma mantle, the entry layer and the exterior cusp or stagnation region, seen in Figure 1.

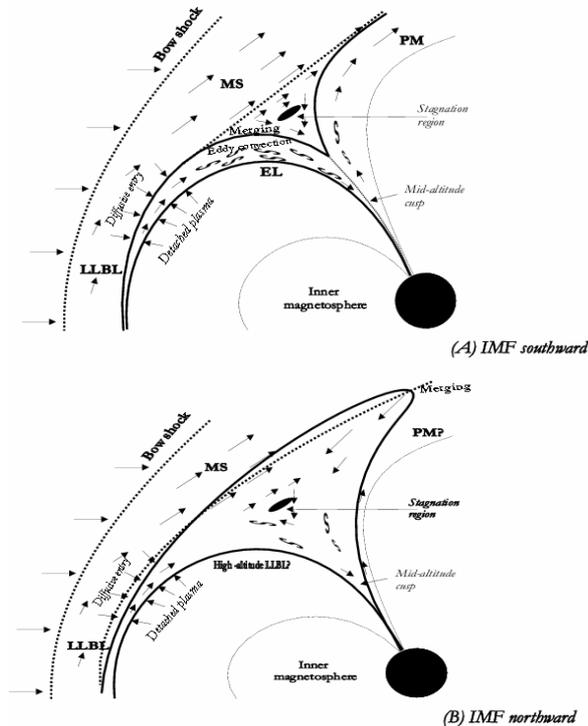


Figure 1. Sketch of the dayside boundary regions related to the polar cusp field lines during southward (after Haerendel, 1978) and northward IMF. MS – magnetosheath; PM – plasma mantle; LLBL – low latitude boundary layer; EL – entry layer.

The plasma mantle is located on the open field lines where the injected magnetosheath plasma continues tailward which was first reported by Rosenbauer et al. [1975]. The plasma density in this region towards sheath density level and the  $\beta \ll 1$ . The entry layer [Paschmann, 1976] is located on the magnetospheric field lines just equatorward of the cusp. It is a region of diffusive, turbulent entry of magnetosheath plasma onto field lines that map to the low-altitude cusp. It has been so termed because it appears to be the region of dominant plasma entry into the magnetosphere. The transport mechanism is likely to be achieved through eddy convection which manifests itself in the irregular,

low speed plasma flow, and may be incited by the turbulence in the adjacent exterior cusp [Haerendel et al., 1978]. The exterior cusp / stagnation region is bounded on the inside by the cusp-like indentation of the magnetopause and outside by the free-flow stream line of the magnetosheath flow which Sckopke et al. [1976, 1981] constitutes a pocket of hot and 'stagnant', possibly turbulent plasma. In fact, as early as 1960's, the stagnation region is already predicted by gas dynamics models [Spreiter et al., 1967, 1980]. Furthermore, this picture has been corroborated by HEOS 2 measurements [Sckopke et al., 1976; 1981]. The stagnation region cannot be linked to the plasma mantle or LLBL in a simple way. A qualitative explanation was given by Haerendel et al. [1978] who noted the similarity of the situation near the cusp to hydrodynamic flow around a corner, in which vortex formation and separation are known to occur and to initiate some level of turbulence (Figure 1). The exterior cusp region appears to be a steady high pressure center of "stagnant" magnetosheath plasma, the flow in this region is rather turbulent, both in magnitude and direction.

The mantle is generally thicker for southward than northward IMF Bz [Sckopke et al. 1976]. These original researchers believed that the plasma mantle is open, and the LLBL is closed. It has been indicated that in the HLBL (Entry layer) the plasma density is almost as high as the magnetosheath but generally lacks the strong antisunward plasma flow. In fact, even sunward flow has been reported by Paschmann [1976]. Lundin et al. [1985] suggested that one of the characteristic features of the entry layer is the strong variability of magnetosheath plasma entry with frequent plasma injections. On the basis of Defense Meteorological Satellite Program (DMSP) F2 data, Newell et al. [1987] indeed observed that the cusp low-altitude latitudinal extent is narrower when Bz is southward than when Bz is northward. This has been interpreted that the enhanced convection flow is too rapid to allow the plasma to reach low altitudes. In non-reconnection models the cusp position and extent are less sensitive to the IMF, but more strongly dependent on the solar wind ram pressure.

Only a few satellites such as HEOS-2, Prognos-7, Hawkeye and Polar have made in-situ observation in the high-latitude boundary layer regions where the proposed entry for the northward IMF takes place. The high latitude boundary layer has scarcely been studied compared to the region around the subsolar point.

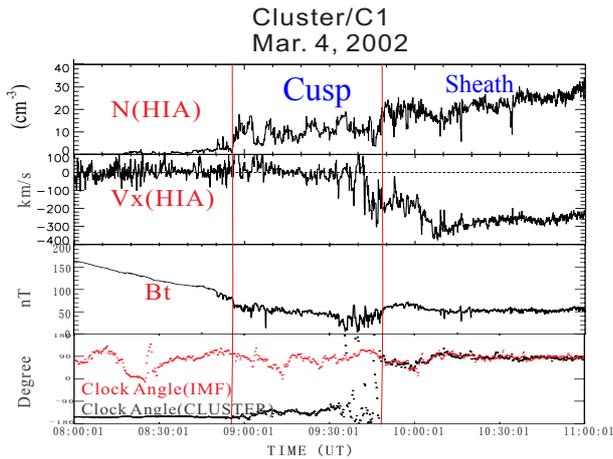


Figure 2: A typical cusp observed by Cluster satellite: CIS plasma moment and magnetic field obtained by Cluster (Rumba) from 08:00 to 11:00 UT, Mar. 4, 2002. The plasma density and plasma velocity  $V_x$  are given in the first and second panels. The magnetic field magnitude (in nT) are shown in the third panel. The magnetic field clock angle obtained by ACE (IMF) and Cluster satellites (local) are plotted in the bottom panel.

### 3. THE DEFINITION OF THE MAGNETOSPHERIC CUSP

In the text book, the polar cusps are usually defined as funnel-shaped areas in the high latitude of both hemispheres with near zero magnetic field magnitude. They provide a direct entry for the magnetosheath plasma into the magnetosphere [e.g., Reiff et al., 1977; Marklund et al., 1990]. However, the definition of the cusp used by MHD simulation [Siscoe et al., 2005] is "a weakening of the magnetic field owing to a pool of magnetosheath plasma within the magnetosphere - since the current is the plasma's diamagnetic current associated with the field weakening" or in concept as "a region of open field lines extending poleward from the open/closed boundary (which is tied to the dayside merging region on the magnetopause) to where particles no longer are able to directly enter". However, there is not always a clear distinction between such a conceptual definition, some observational identifications, and actual determinations of when the cusp is really being observed by a given spacecraft. The primary and the most widely used method of identifying the cusp is by means of a combination of plasma and magnetic field observations, although just plasma or magnetic field measurements have been used in the past in some cases. Using both sets of observational data the cusp has been defined as a high latitude region with a population of particles of shocked solar wind energies and density somewhere within or near the local noon sector, the criteria are:

- I. turbulent and depressed magnetic fields,
- II. High density plasma ( $\sim$  sheath level),
- III. Stagnant plasma flow ( $V_x \sim 0$ ),
- IV. The clock angle criterion (the Cusp clock angle should be different with the IMF's), an example of observed cusp is shown in Figure 2.

A number of questions that now need to be investigated in detail: What is the importance of the cusp to the physics of the magnetosphere and the topology of the front side high-latitude magnetopause? What's the nature of the boundaries between different regions? What's the plasma transport mechanism through the cusp and the boundary layers? Are the observed double or triple cusps temporal or spatial effect? How are they formed? What's the role of the cusps in supplying plasma to the plasma sheet? Are there multiple cusps active at a given time? Or does a single region move around more and faster than its low altitude counterpart? Only with measurements at larger separations and more coordination of multiple satellite measurements for particular cusp crossings will it become evident what the true nature of the cusp is and what roles the cusps play. Understanding the polar cusps is essential for a thorough understanding of the entire physics of the magnetosphere, and of the dynamical interaction between the solar wind and any planetary magnetosphere.

### 4. ENERGETIC PARTICLES IN THE HIGH LATITUDE BOUNDARY/CUSP REGION

#### 4.1 OBSERVATIONS

Figure 3 shows the different regions -- southern HLTR/cusp, radiation belt and northern HLTR/cusp as obtained by Cluster/RAPID during these two consecutive orbits. The geomagnetic index Dst is shown in the top panel. The first two cusp crossings (1 and 2) happened during a rather quiet time period; the Dst indices were small and positive. The later two crossings (3 and 4) occurred in the main phase of a strong magnetic storm. As we can see from Figure 3, the high latitude magnetosphere (both northern and southern HLTR/cusp) regions are two of three locations that energetic particle are encountered. Both southern and northern high latitude boundary and/or cusp regions can be distinguished easily by the fact that (1) the ion flux increased sharply in all energy channels from 30 to 400 keV, whereas (2) the energetic electron flux which appeared during quiet time was not present during disturbed time.

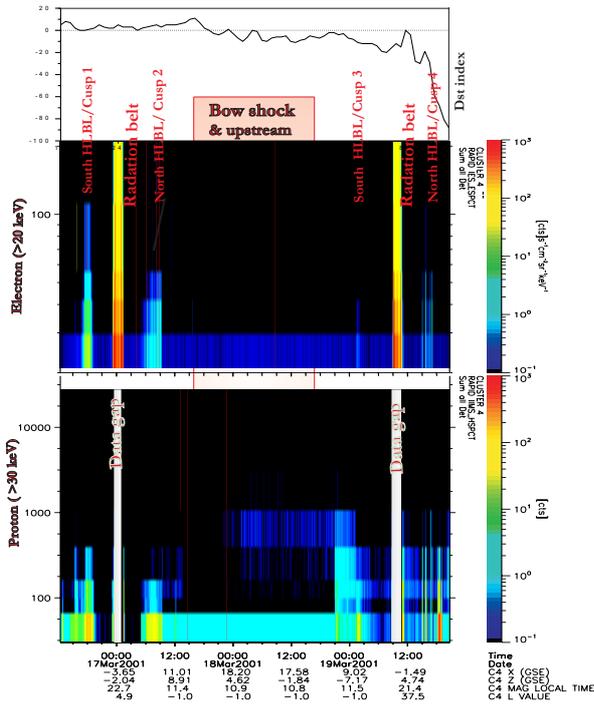


Figure 3: An overview of RAPID data from 09:00 UT, March 16 to 21:00 UT, March 19, 2001 together with geomagnetic activity Dst index. From the top the panels show: Dst index; electron spectra from 20 to 400 keV; and proton spectra from 30 to 2000 keV. The marks indicate the different regions -- the southern HLBL/cusp, the radiation belt and the northern HLTR/cusp which Cluster experienced during a full orbit. The 'cusp' here refers all the high latitude magnetospheric regions, the 'radiation belt' here refers all inner magnetospheric regions, and two data gaps are indicated [Zong et al., 2005].

During the above two quiet time high latitude boundary/cusp crossings, there were pronounced fluxes of electrons in the high latitude boundary/cusp region, indicating either a closed field line geometry in the cusp region, a special open field line configuration that could trap electrons very efficiently for a long time or a long-lived source supplying these electrons to open field lines. These electrons lasted about 2 hours (from 1706 to 1907 UT, 16 March 2001) and 3 hours 20 min (from 0600 to 0920 UT, 17 March 2001), respectively. Furthermore, no obvious substorm injections were observed by the Los Alamos satellites for both of the above quiet time high latitude boundary/cusp crossings. Further, there were no energetic electron events observed by ACE in the upstream interplanetary space during 16-19 March. Thus, these observed electrons should not be solar energetic electrons as described by

Lin [1985] and Klassen et al. [2002]. The lack of substorm activity and high fluxes of electrons upstream at ACE indicates that the observed electrons are locally trapped electrons rather than substorm injected electrons drifting to the high latitude region or solar flare related electrons. In contrast, two consecutive HLBL/cusp crossings during geomagnetically disturbed times are.

## 4.2 SINGLE PARTICLE SIMULATION

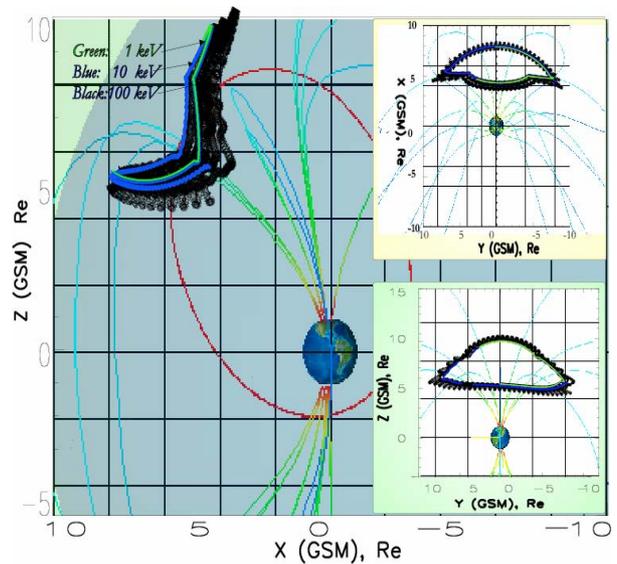


Figure 4: The trajectory of protons (1, 10 and 100 keV) with  $90^\circ$  pitch angle in the Tsyganenko 96 model.

We have seen in session 4.1 there are energetic particles in the high latitude boundary/cusp region during northward IMF. Statistical study showed that energetic ions have been observed in the cusp region during most of the crossings (80%) [Zhang et al., 2005]. This indicates that those particles are trapped in the high latitude region. In this session, we will interpret this observation with single particle simulation.

According to the traditional dipole field model, the dayside high latitude or cusp region cannot trap particles [Roederer, 1970; 1977]. The cusp region of the ideal dipole field is not an "excluded region" in the Störmer theory [Störmer, 1911]. This means that, in the high latitude region, the particles cannot be trapped for much longer than the bounce time; the  $E \times B$  drift will take the particles away. However, the dipole field can be modified fundamentally by interaction with the solar wind. The outer cusp regions where the magnetic field lines either close in the dayside sector or extend into the night side sector over the polar cap could be caused by

reconnection. This region is a region of weak magnetic field, which is directly from the interaction of the solar wind with the geomagnetic field predicted by Chapman and Ferraro [1931] by using a simple image dipole, as shown by the magnetic field model of Antonova and Shabansky [1968].

Instead of a dipolar field, the cusp region appears to be quadrupolar. However, the importance of the existence of an off-equatorial B-minimum in the outer cusp has been underestimated for a long time, although it could be of extreme importance for understanding the behavior of energetic particles in the magnetosphere. Antonova and Shabansky [1968] and Shabansky [1968] noted that, with a minimum magnetic field existing off the equator in the outer cusp region, charged particles would not drift but rather branch off towards the magnetic field minimum at high latitudes. This has also been supported by in situ magnetic field measurements [Zhou et al., 1997].

Shabansky [1971], Antonova and Shabansky [1975] provided observational evidence for the trapping of energetic particles (of several tens of keV, up to a few hundreds of keV) in the high latitude region. Sheldon et al. [1998] pointed out that an energetic electron will drift on a closed path around the front of the magnetosphere, and found that electrons could be trapped in the outer cusp. In fact, a temporary trapping in the cusp field minimum was first examined by Delcourt et al. [1992]. Further, Delcourt and Sauvaud [1998, 1999] pointed out that, under the effect of the cusplike mirror force near the dayside magnetopause, energetic plasma sheet particles initially mirroring near the equator are expelled from low latitudes and subsequently swept into the boundary layer at high latitudes.

Figure 4 shows trajectories of test protons (1, 10, 100 keV) launched with  $90^\circ$  pitch angle from the cusp region. The trajectory tracing was performed using the Tsyganenko 96 model. In this calculation, the full particle dynamics have been considered, not just the guiding center computation; the calculation was performed using a fourth-order Runge-Kutta technique with a time step adjusted to some fraction of the particle gyration periods. It can be seen from Figure 8 that the test protons launched from the local minimum magnetic field region encircle the outer cusp region; all of the protons experience a pronounced bouncing motion in the high latitude region which differs from mirroring motion on either side of the equator (as ring current ions). Figure 4 shows that the ion trajectories in the outer cusp region are somewhat similar to those on L shells of a dipolar magnetic field. The limiting second invariant of these trapped orbits occurs when the mirror

point  $B_{min}$  approaches the dayside equatorial field strength; in the local gradient field they drift away from the cusp.

These ion trajectories exist both on the dayside (equatorward, with closed magnetic field lines) and the mantle region (poleward, with open magnetic field lines). This behaviour follows from the existence of a local B minimum during the drift path from the closed field lines region to the open field line region in the frontside magnetosphere.

It should be pointed out that large-scale magnetospheric convection is not accounted for in the present modelling results in the trajectory computation. If convection is accounted for, as pointed out by Delcourt and Sauvaud [1999], these closed drift paths in the outer cusp may be opened (see Figure 13 of Delcourt and Sauvaud [1999]). However, the present modelling results could apply to quiet times when magnetospheric convection is reduced and the convection electric field may be only 5% of its value during active times.

The observed energetic electrons in the high latitude boundary region may be provided by tail plasma sheet particles because of a minimum magnetic field existing off equator in the high latitude region of the magnetosphere. Delcourt and Sauvaud [1998, 1999] pointed out that under the effect of the cusplike mirror force near the dayside magnetopause, energetic plasma sheet particles initially mirroring near the equator are expelled from low latitudes and subsequently swept into the boundary layer at high latitudes. Both electrons and ions can be stably trapped in the high latitude region during quiet periods. This conclusion is supported by both the observations and the modelling results mentioned above. As magnetospheric convection is enhanced, the electrons initially trapped in the high latitude region could be de-trapped [Delcourt and Sauvaud, 1998, 1999]. In fact, no stable trapped electrons were observed during active times (see Figure 3). These de-trapped electrons could further form an electron layer just outside the magnetopause as observed [e.g., Meng and Anderson, 1970 and Baker and Stone, 1977].

## 5. THE BOUNDARY BETWEEN THE CUSP AND MAGNETOSHEATH

The high latitude boundaries include boundary between the magnetosheath and cusp, the boundary between the magnetosheath and the High Latitude Trapping Region (HLTR) which is the closed field line region on the dayside in the high latitude region, the boundary between cusp and HLTR and the boundary between

mantle and cusp. The properties of the high latitude boundaries vary rather dramatically under different solar wind conditions. We present statistical results based on 4 years of data obtained by Cluster when these spacecraft were in the vicinity of the dayside magnetopause. During northward Interplanetary Magnetic Field (IMF), the interfaces between the magnetosheath and cusp are rather clear. The changes of the energetic particle flux, plasma temperature, density and velocity across the magnetopause under northward IMF were analyzed by superposed epoch analysis. The plasma flow and density decrease and the proton temperature increases across the magnetopause from the magnetosheath into the cusp. Further, during extreme storm times, the cusp is more turbulent than during quiet times and there is no clear plasma density change across the magnetopause.

### 5.1 BOUNDARY AND CLOCK ANGLE OF THE IMF

The boundary between the magnetosheath and the cusp has been studied by Lavraud et al. We have found that the boundary between the magnetosheath and the cusp is clear sometimes but unclear at other times. The definition for a clear boundary is: There is a jump in plasma flow ( $> 30\text{ km/s}$ ) (e.g., a flow change from  $100\text{ km/s}$  to  $0$ ) and at least two components of the magnetic field ( $> 5\text{ nT}$ ) (e.g., the magnetic  $B_y$  and  $B_z$  change from  $-10\text{ nT}$  to  $+10\text{ nT}$ ).

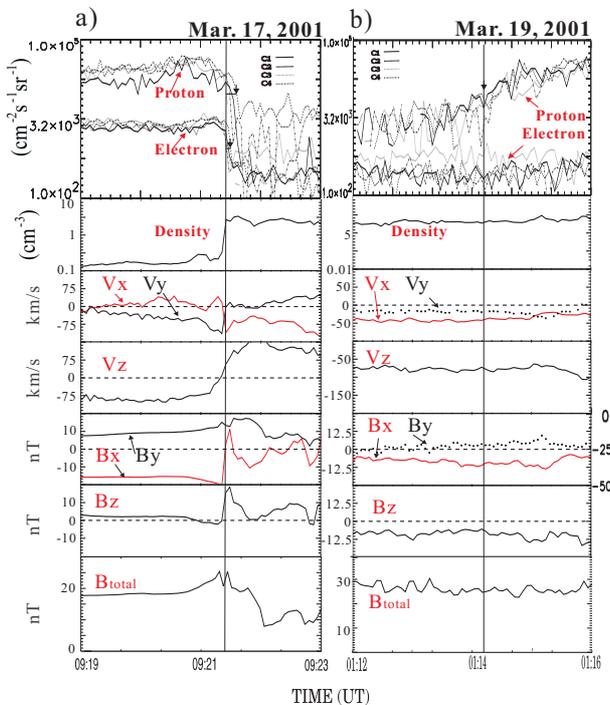


Figure 5: Examples for clear and unclear boundary. a) Clear boundary on Mar. 17, 2001 b) Unclear boundary on Mar. 19, 2001. From the top, the panels indicate

integral electron and proton fluxes; plasma ion density, plasma ion velocity  $V_x$ ,  $V_y$  and  $V_z$  (in km/s), and superposed magnetic field  $B_x$ ,  $B_y$  and  $B_z$  components for spacecraft 1 (Rumba).

Figure 5 shows an example for a clear and unclear boundary. Figure 5 a) shows a clear boundary observed by Cluster when it travels outbound from the northern magnetosphere into the magnetosheath on March 17, 2001. From this figure we can see that all the parameters including energetic proton and electron flux, plasma density, velocity and magnetic field have clear boundaries. Figure 5 b) shows an unclear boundary observed by Cluster when it travels inbound from the magnetosheath into the southern magnetosphere on March 19, 2001. We can see from this figure that all the parameters except energetic proton flux have no clear boundaries. We don't know exactly where the boundary is since the parameters change smoothly from the magnetosheath in the magnetosphere if we look at the plot for a longer time which is not shown here, so we put the vertical line which indicates the boundary location at the location where energetic proton flux changes a lot.

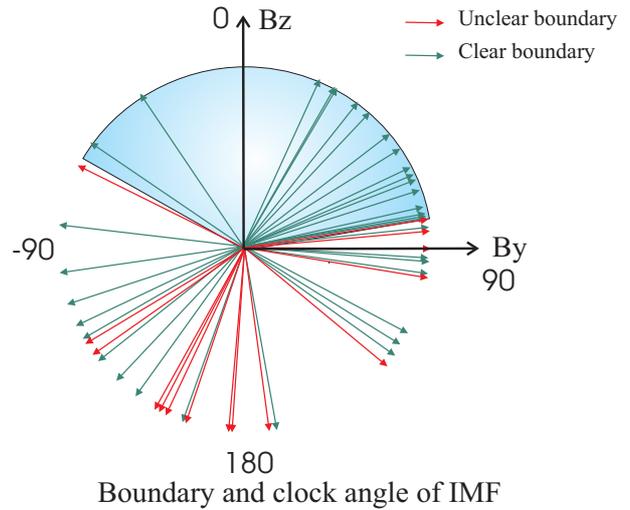


Figure 6: The IMF clock angle dependence of the boundary between the magnetosheath and cusp.

We have surveyed Cluster data in 2001 and 2002 and found that almost all the boundaries between the HLTR and the magnetosheath are clear (looks like Figure 5 a)). However the boundary between the cusp and magnetosheath is more complicated. Sometimes it's clear and sometimes it's unclear (looks like 5 a) and b) respectively). It's found that there is some relationship between the boundary and the clock angle of the IMF. In Figure 6, all the boundaries between the cusp and the magnetosheath in 2001 and 2002 are shown. The arrows

indicate the IMF direction projected in the GSE YZ plane. The red arrows indicate unclear boundaries and green arrows indicate clear boundaries. In the shaded region, all the arrows are green which means when the IMF clock angle is between  $-65$  and  $81$  degree, the boundary between the magnetosheath and the cusp is clear.

## 5.2 QUIET TIME VS. EXTREME STORM TIME

The properties of the high latitude boundaries vary rather dramatically under different solar wind conditions. In order to study the average variations of key plasma parameters in the vicinity of the magnetopause under different conditions, we perform a superposed epoch analysis. We present statistical results based on 4 years of data obtained by Cluster when these spacecraft were in the vicinity of the dayside magnetopause.

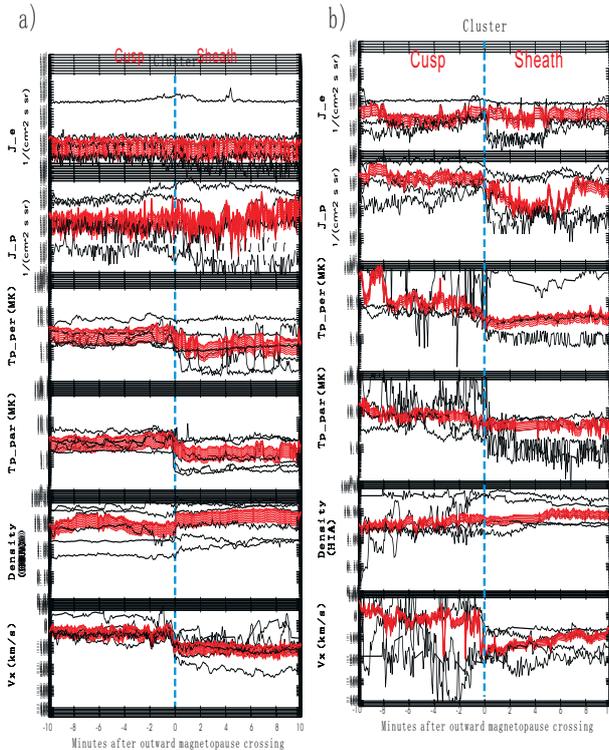


Figure 7: Superposed epoch analysis of the energetic particle flux, the plasma temperature, density, and velocity change from cusp region across the magnetopause a) under northward IMF; b) during extreme storm time ( $Dst < -100$  nT).

Figure 7a) shows superposed epoch analysis of the energetic particle flux, the plasma temperature, density, and velocity change from cusp region across the magnetopause under northward IMF in northern hemisphere. The vertical dashed line marked the

magnetopause position which is identified by the jump in plasma parameters including temperature, density and velocity. The x axis is the minutes after outward magnetopause crossing. The time interval in this plot is 20 minutes, 10 minutes before and 10 minutes after the magnetopause crossing. During northward IMF, the interfaces between the magnetosheath and the cusp are rather clear. The plasma flow and density increase and the proton temperature decreases across the magnetopause from the cusp into the magnetosheath.

Figure 7 b) shows superposed epoch analysis of the same parameters as a) but during extreme storm time ( $Dst < -100$  nT). By saying a event is during extreme storm time we means that the most negative  $Dst$  during one storm is  $< -100$  nT and the event is observed during the storm time (initial phase, main phase or recovery phase). In Figure 7 b) all the events during extreme storm time from 2001 to 2004 are included. If the magnetopause crossing is in the southern hemisphere, we reverse the time sequence so that the crossing is still from the magnetosphere into the magnetosheath. Compare Figure 7 b) and Figure 7 a), we can see that during extreme storm times, the cusp is more turbulent than during quiet times. We also noted that there is no clear plasma density change across the magnetopause during extreme storm time.

## 6. THE CUSP DYNAMICS

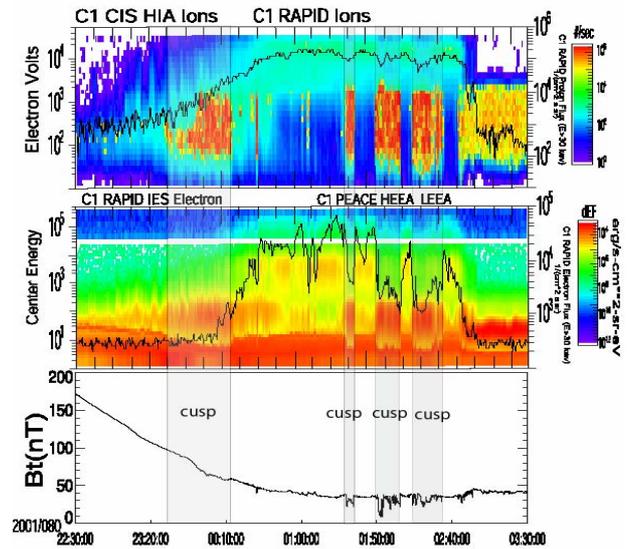


Figure 9: Four-fold cusp observed by CLUSTER. Plasma ion and electron spectra are over-plotted with energetic particle in the first two panels. Last panel shows the total magnetic field.

Unraveling the structure of the dayside cusp is a major CLUSTER objective. Four cusp-like regions were observed consecutively in about five and half hours on March 21–22, 2001 by all four CLUSTER spacecraft

when the IMF was northward with a significant  $B_y$  component. The four-fold cusp (Figure 9) was surrounded by the dayside magnetosphere rather than the magnetopause. Evidence was presented indicating that the multiple cusps were probably a temporal sequence. The boundary normal, velocity and timing analysis for six clear boundaries of the cusps indicated that the observed cusps encounters are mostly caused by oscillation of the single cusp which was shifted back and forth between the dayside magnetosphere/trapping region and the cusp region. The normal velocities at boundary interfaces for exits from the cusp were found to be almost three times as large as that for entry into the cusp.

These observations suggest that the shape and location of the cusp is often changing as a result of dynamic processes in the high latitude regions. Furthermore, by combining the four Cluster spacecraft positions and crossing times at the interfaces, we are able to determine the normal velocity and direction of the discontinuity based on the triangulation method [Russell et al., 1983]. Assuming planar discontinuity and the speed of the discontinuity was constant in time and space over the Cluster separation distance, the equation is simply given by

$$Rl \cdot n = Vn \cdot tl$$

Here  $Vn$  is the normal velocity of the discontinuity,  $Rl = (r12, r13, r14)$  is a tensor with  $r\{ij\}$  being the spacecraft separation vectors and  $tl = (t\{12\}, t\{13\}, t\{14\})$  consists of the differences between the crossing times of the corresponding spacecraft. The obtained normal speeds for the six interfaces are given in Figure 10 (between panel 2 and 3). As we can see from panel 3 and 4, the polar  $\theta$  angles shows little change in the first 4 interfaces, however, they change signs for the last two interfaces. The wavelike motion can be more clearly seen in the  $\phi$  angle, the direction of exiting the cusp at 2, 4, 6 is opposite to the entering direction based on the azimuthal angle.

Although Cluster observed four cusp regions, the analysis shows that three of these cusps (later three cusps), which are located in the lower-latitude, are a temporal feature which can be attributed to cusp boundary movements or possibly wave activities.

During this event, the IMF was steady with IMF  $B_z > 0$ , and IMF  $B_y$  is slightly larger or comparable with IMF  $B_z$ . Between 0100 and 0330 UT on March 22, 2001, the solar wind speed was 300 km/s and the radial dynamic pressure was around 3 nPa obtained by Wind satellite located at GSM (-11.7, -183.6, -109.6) Re, see Figure 11. The time lag between Cluster and Wind is about 6 min

with Cluster in advance. A solar wind pressure pulse encountered the Earth at ~0058 UT, Mar.22, 2001, shortly before the second cusp was observed by Cluster while there was a change of solar wind azimuthal flow (shown in Figure 11). The solar wind East/West flow changed from 6 degree to about 2 degree, and the  $V_x$  change from -280 km/s to -310 km/s. The Earth's magnetosphere bears an analogy to the windsock response to changes of the solar wind component flow as suggested by Zong et al. [2004]. When the solar wind azimuthal flow encountered the Earth, the ratio of the Y and Z components of the solar wind dynamic pressure to solar wind thermal pressure  $P_{dy}/P_{th}$  became 29%. Therefore, the position of the cusp will be changed. Cluster spacecraft entered the cusp, then the solar wind dynamic pressure or other effects shift the cusp back and forth three times as if Cluster flew through three cusps.

On the other hand the second and third observed cusps don't seem to have clear links with the solar dynamic pressure or flow components. Nevertheless, the mentioned solar wind pressure pulse could trigger the magnetospheric boundary wave [Sibeck et al., 1998].

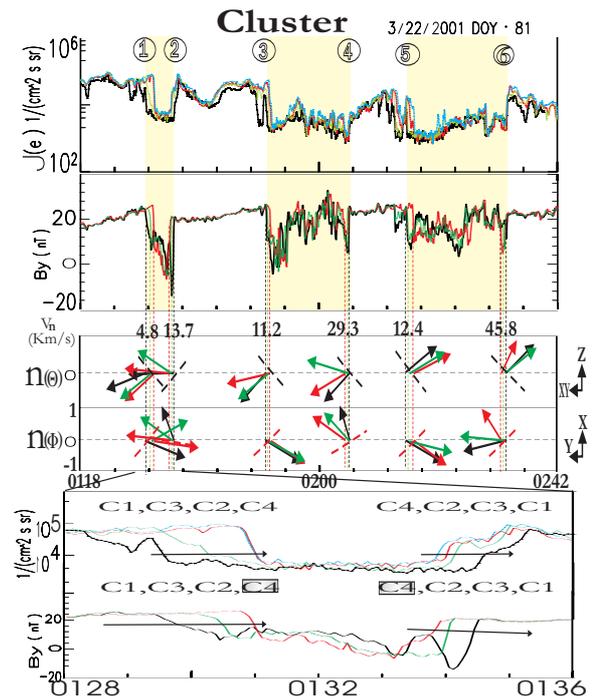


Figure 10: Electron flux, Magnetic field  $B_y$  component measured by the four Cluster spacecraft during the time period from 01:00 to 03:30UT on Mar 22, 2001. The boundary normals were determined by MVA for all available spacecraft during cusp crossings. One of the three cusps is expanded to see the spacecraft crossing order more clearly.

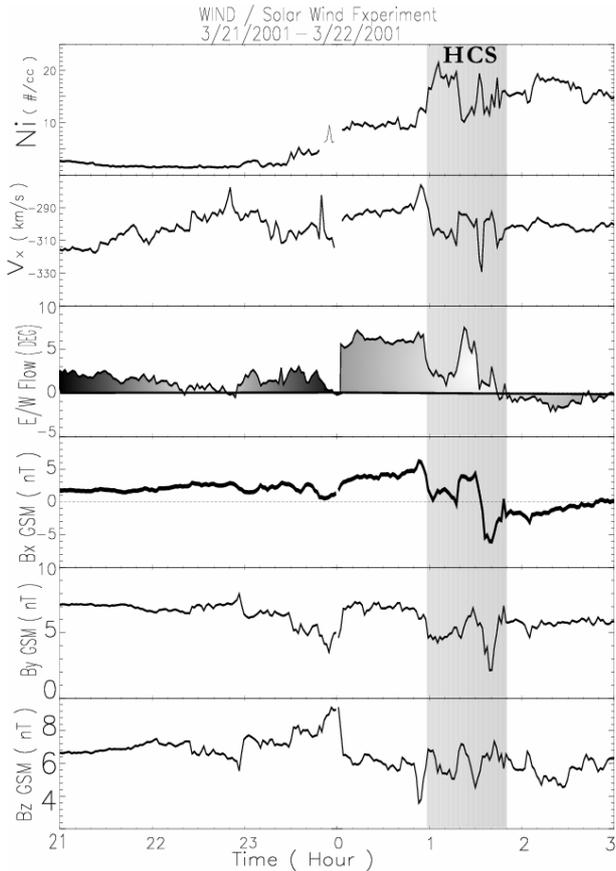


Figure 11: Solar wind and IMF observed by the Wind satellite. The Heliospheric Current Sheet (HCS) which is identified by the high plasma density and sign change of the IMF  $B_x$  is marked. The time of this Figure is shifted according to the solar wind velocity.

When the IMF is northward, a cold and dense plasma sheet is often observed [baumjohann et al., 1989, Fujimoto et al., 1996, Fujimoto et al., 1998, Terasawa et al., 1997]. The magnetosheath plasma near the cusp region is relative cold comparing with magnetospheric plasma, dense and almost stagnant. When cusp reconnection occurs, the newly reconnected flux tubes tailward of the cusps in both hemispheres sink and contract into the magnetosphere. Subsequently it sweeps around the flank, and is convected tailward. As the plasma is captured and transported to the tail, it is moderately heated near the reconnection site, and the temperature is just below 1 keV [Song et al., 1992]. The density in the captured flux tubes is characteristic of cold, dense and almost stagnant plasma. The values of temperature, density and low flow speed for the captured plasma are in good agreement with Geotail observations in the dusk flank as shown in Figure 12.

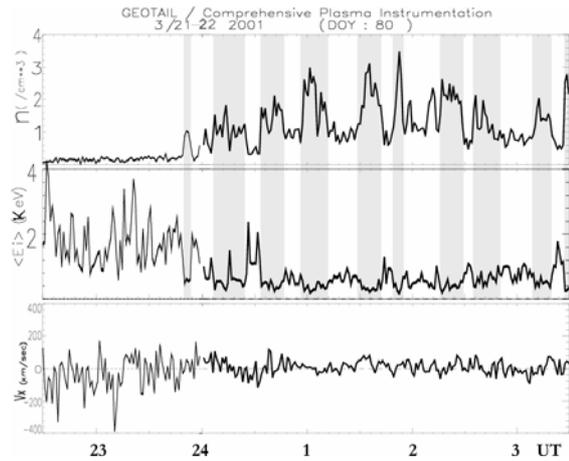


Figure 12: Geotail observations at the duskside tail-flank. From top, the ion density, average energy (temperature, in keV), and the ion bulk flow  $V_x$  in km/s. bursty-like cold-dense plasma are observed.

It should be noted, however, the cold dense plasma sheet is also waving. The period is about 20 min. This agrees with the cusp oscillating period (a rough period of 22 min.) very well. This is an additional evidence to support the idea that the cold dense plasma sheet observed by Geotail is closely related the oscillating cusps which are probably formed by high latitude reconnection during the extended northward IMF period.

## 7. CONCLUSIONS

Energetic ions and electrons have been observed during quiet time (IMF  $B_z$  has a predominate positive component). The energetic electrons disappeared during storm time. Single particle simulation shows that energetic ion could be temporally trapped in the High Latitude/Cusp Region whereas electron could not be.

When IMF is northward, the interfaces between the magnetosheath and magnetosphere are rather clear. However, this interface will become uncertain when IMF turns southward. The plasma density decrease and the proton temperature increases across the magnetopause from the magnetosheath into the cusp. During extreme storm time, the cusp is more turbulent than quiet time and there is no clear plasma density change across the magnetopause.

Multiple Cusps are observed by Cluster satellite in the high latitude region, in the meantime, the cold-dense plasma with fluctuations are observed in the dusk-side of the tail plasma sheet by Geotail satellite. This is consistent with the idea that the high latitude reconnection during northward IMF is the responsible

mechanism of the formation of the cold-dense plasma sheet. The observed multiple cusps are very possibly temporal sequence. The cusp was shifted by the solar azimuthal dynamic pressure or wave back and forth three times in about 5 and half hours interval as if Cluster flew through the cusp four times. Further we suggest that the solar wind azimuthal flow is the controlling factor of the cusp position and is as strong as, potentially even stronger than, that of the IMF  $B_y/B_z$  component. The importance of the solar wind azimuthal and north/south flow as a dynamic driver of the cusp, and even the whole magnetosphere has been more or less neglected or underestimated.

The full impact of the cusp is going to be evident when we have multiple satellites displaced from one another by large distances (from a large fraction of an Earth's radius to a few Earth radii), as well as satellites located within 100 km of one another and observing using interferometric techniques. New Cluster orbits will make possible study of the nature of particle boundaries within the cusp and high latitude regions to resolve the mechanisms that transport, and possibly accelerate, the thermalized plasma and energetic particles in and through the cusp and boundary layers. The multi-scale Cluster separations will help to resolve the issue of whether the observed multiple cusps are a temporal or spatial effect, as well as determining the size of the cusp. The question of how multiple cusps are formed will be addressed. An outstanding question is the role and importance of the cusp in supplying plasma to the plasma sheet and the relationship between high latitude reconnection and the appearance of cold dense plasma within the plasma sheet.

### Acknowledgements

A part of this work by Fu and Pu at PKU is supported by CNSF Grant 40390152, 40425004, Chinese Fundamental Research ProjectG200000784,

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