Structure of the magnetopause boundary layers discovered by Cluster multipoint observations

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(Cluster-Double Star symposium, Sep. 19-23, 2005)

Outline

Detection of Rolled-up Kelvin-Helmholtz Vortices Cluster + Realistic three-dimensional (3D) MHD simulations

Two-Dimensional Structure of the Magnetopause and Flux Transfer Event (FTE) Cluster + Grad-Shafranov reconstruction technique

Model-based data analysis helps us a lot to interpret data obtained by the Cluster multi-spacecraft measurements.

Why Kelvin-Helmholtz vortices important?

• Vortices, developed through the nonlinear growth of the Kelvin-Helmholtz instability (KHI) at the flank magnetopause, can be the agent of transport of solar wind plasmas into the magnetosphere under northward IMF conditions, which has been a long-standing problem.



Plasma transport processes accompanied by the KHI growth

Magnetic reconnection within a rolled-up KH vortex (e.g., Otto & Fairfield, 2000)

• Collapse of vortices mediated by electron inertia effects (Nakamura et al., 2004)

• Turbulence triggered through Rayleigh-Taylor instability in a rolled-up vortex (Matsumoto & Hoshino, 2004)

Numerical simulations suggest that all these transport mechanisms can occur ONLY when the KHI has grown to form "Rolled-up" vortices.

Can "rolled-up" KH vortices form in a tail flank-like situation?



Simulation results suggest that the roll-up of KH waves can be achieved as long as the PS is thick enough (the thickness is comparable to, or larger than, the KHI wave length).



(Takagi et al., 2005)

To what extent does the KHI grow in the actual magnetosphere?



Single- (or dual-) spacecraft measurements (e.g., Kivelson & Chen, 1995) could not answer this question.

Multipoint measurements by the four Cluster spacecraft can answer.





Evidence of plasma transport across the magnetopause Log(Energy flux) (keV s⁻¹ cm⁻² sr⁻¹keV⁻¹) (ла) 10 10 10 10 6 5 10' 600 € 400 Low energy ions of م 200 sheath origin 0 5 0 0 5 0 C3 C41 detected throughout Log(Density) (cm⁻³) -5000 0.8 C1 Y (km) 0.6 0.4 5000 0.2 C3 10000 Sat 1 CLUSTER HIA H+ (Product 23) .USTER HIA H+ (Product 2B) Sat1CL 20:30 20:32 34 2001-11-20/20:32:08->20:32:20 11-20/20:34:33->20:34:45 200 1.07 10⁻⁶ 16 1.88 10⁻⁶ 7.62 5.58 1500 1.45 10⁻⁷ 2.55 10⁻⁷ 1000 1000 1.97 10⁻⁸ 3.46 10⁻⁸ Perp (ExB) (km/sec) Perp (ExB) (km/sec) cm^3) 500 500 2.68 10⁻⁹ 4.70 10⁻⁹ 3.63 10⁻¹⁰ 6.39 10⁻¹ 0 0 4.93 10⁻¹¹ 8.68 10 -500 -500 6.68 10⁻¹² 1.18 10⁻¹¹ distribu -1000 -1000 1.60 10⁻¹² 9.07 10⁻¹³ .23 10⁻¹³ dist -1500 -1500 -1000 -500 -1500 -1000 -500 0 500 1000 1500 500 1000 1500 0 V Para (km/sec) V Para (km/sec)

The observation is consistent with transport via KHI!

1-SC detection of "Rolled-up" vortices possible? Difference between Rolled-up & Not rolled-up vortices



Vx vs N seen in simulated data







Comparison with vortices observation

MHD simulation

Cluster observations of rolled-up vortices





When SC separation of Cluster is small,

- Detection of a parent rolled-up vortex can be made by either of the four spacecraft, by identifying Low-density & High-speed flows.
- Then,

• The nature of small-scale waves excited, or thin current sheets formed, in the vortex can be investigated in detail with the help of the multi-point measurements.



Vol 436 11 August 2005 doi:10.1038/nature03931

nature

Small-scale vortices discovered in the CUSP (Sundkvist, et al., 11 Aug. 2005 issue of Nature)

In situ multi-satellite detection of coherent vortices as a manifestation of Alfvénic turbulence

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2D structure of the magnetopause and FTE

Grad-Shafranov reconstruction technique

(e.g., Hau & Sonnerup, 1999) A spatial initial value problem

Assumptions

The plasma structure is:

$$\rho \overrightarrow{\partial t} + \rho (\overrightarrow{V} \nabla) \overrightarrow{V} = \overrightarrow{J} \times \overrightarrow{B} - \nabla \cdot P$$

$$\vec{j} \times \vec{B} = \nabla p$$

- in a magnetohydrostatic equilibrium (time-independent).
- 2-D (invariant along some direction, z)

Grad-Shafranov (GS) equation (e.g., Sturrock, 1994)

$$\frac{\partial^2 A}{\partial x^2} + \frac{\partial^2 A}{\partial y^2} = -\mu_0 \frac{dP_t}{dA} = -\mu_0 j_z(A) \qquad (A = A_z)$$
$$\vec{B} = (\partial A/\partial y, -\partial A/\partial x, B_z(A)), \qquad P_t = (p + B_z^2/2\mu_0)$$

A magnetic field map, A(x,y), is constructed from explicit integration of the GS equation, using measured magnetic fields as spatial initial values.

FTE reconstruction \rightarrow Verification of flux rope models



- Flux rope size ~ 1 Re
 Strong core field
 → Evidence of
 "component" merging
 No reconnection activity any more
 Moving poloword
 - Moving poleward

Total transverse magnetic flux within the flux rope = 0.0549 T·m

- Reconnection *E* field
- (total magnetic flux) /
 - (FTE occurrence period)
- = 0.0549 (T·m) / 5 (min.)
- = 0.183 (mV/m)
- \sim reconnection rate = 0.04

Reconstruction of a magnetopause on 5 July, 2001

A future possibility

• There is a Grad-Shafranov-type equation to describe stream lines.

Hopefully, it might become possible to reconstruct a 2D map of the flow velocity field, for example in KH vortices, from SC

measurements.

Summary

A combination of the Cluster multipoint observations and numerical simulations has enabled us to unambiguously detect "rolled-up" KH vortices at the flank magnetopause.

The detection of the rolled-up vortices is now possible with a 1-SC data, providing a possibility of studying coupling between the vortices and small-scale processes with the Cluster data.

Grad-Shafranov reconstruction of magnetopause and FTE structures using the Cluster data demonstrates time evolution of the magnetopause structure, and provides information on the appropriate FTE models, and on the nature of magnetopause reconnection (the reconnection rate, orientation of X-line, component merging, etc.).

How can the secondary velocity shear be produced?

$$\frac{M_1 v_1^2}{r} = \frac{M_2 v_2^2}{r} = a \quad leads \ to, \qquad \mathbf{V2}$$

$$\begin{cases} v_1 = \frac{a\sqrt{r}}{\sqrt{M_1}} \\ v_2 = \frac{a\sqrt{r}}{\sqrt{M_2}} \end{cases} \qquad \mathbf{M1} \quad \mathbf{V1} \\ \mathbf{M1} \quad \mathbf{VI} \end{cases}$$

$$\Delta V_{SE}$$

$$\therefore |v_1 - v_2| = a\sqrt{r} \left| \frac{1}{\sqrt{M_1}} - \frac{1}{\sqrt{M_2}} \right| = \frac{a\sqrt{r}}{\sqrt{M_1}} \left| 1 - \sqrt{\frac{M_1}{M_2}} \right|$$

$$\mathsf{M2}$$

$$\Delta V_{SEC} \propto \sqrt{r_{SEC}} \left(1 - \sqrt{\frac{1}{N_{RATIO}}} \right)$$

sec = curvature radius of he 2nd velocity shear layer

At a certain radial distance from the vortex center, the centrifugal force exerting on the low-density and dense fluids must be equal.

Then, the shear velocity depends on the mass ratio and on the curvature radius of the interface between the two fluids.

Hall (two-fluid) MHD equations including electron inertia effects

 $n\frac{d\vec{V_i}}{dt} = -\nabla P + \vec{J} \times \vec{B}$

Momentum equation

Equation of state

$$\frac{\partial}{\partial t} \left(1 - \frac{1}{M} \Delta \right) \vec{B} = \nabla \times \left[\vec{V_e} \times \left(1 - \frac{1}{M} \Delta \right) \vec{B} \right]$$

$$\lambda_e = \sqrt{1/M} = \sqrt{m_e/m_i}$$

Induction equation including finite electron inertia

$$\vec{V_e} = \vec{V_i} - \vec{J}/n$$

$$\vec{J} = \nabla \times \vec{B}$$

