

# MAGNETOTAIL SUBSTORM FEATURES FROM MULTI-POINT OBSERVATIONS

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## ABSTRACT

We examine several dipolarizations and a reversal of plasma flow detected by Cluster at the downstream distance of  $\sim 19 R_E$  in the midnight sector of the magnetotail on August 22, 2001. These substorm features occurred during a sequence of substorm activity seen by the IMAGE/FUV imager. Findings consistent with previous reports are the following: (1) there is generally a decrease of the  $z$ -component of the magnetic field ( $B_z$ ) prior to dipolarization; (2) the onset of dipolarization can precede the onset of significant plasma flow; (3) dipolarization is quite episodic in nature, consisting of multiple increases and decreases of the  $B_z$  component. New features revealed in this study are the following: (1) the turbulent region associated with dipolarization can exhibit dynamo effect as well as dissipation, (2) the  $x$ -component of the Lorentz force ( $j \times B$ ) <sub>$x$</sub>  does not change sign at plasma flow reversal although the other two components do (indicative of changes in the magnetic stress consistent with shifting dominance of multiple activity sites rather than a tailward movement of a single activity site as commonly assumed), and (3) the frozen-in condition breaks down during dipolarization in several regions where the  $B_z$  component and the total magnetic field strength have large values, different from the expectation of the diffusion region associated with magnetic reconnection.

## 1. INTRODUCTION

Plasma environment in the Earth's magnetotail has been studied since the early days of the Vela satellites in the late 1960's (see, e.g., [1]). Plasma flows were inferred from these early measurements even though the plasma detectors and the satellite spin axis orientation were not ideally designed to obtain accurate determination of plasma flow direction. It has been recognized in early substorm research that high-speed plasma flows occur

during geomagnetic active periods. The lack of multi-point observations in the magnetotail and the ambiguity in differentiating temporal from spatial variations have posed major obstacles in advancing our understanding of substorm phenomena in the magnetotail. In spite of these deficiencies, attempts have been made to construct global substorm development in the magnetotail from single-point observations through statistical studies [2-5]. Synoptic maps of substorm features from these studies have enabled substorm models incorporating disturbances in the magnetotail [6-7].

The ISEE-1/2 mission demonstrated the power of multi-point measurements from satellites flying in close formation. Subsequently, the International Solar Terrestrial Physics (ISTP) program has brought the prowess of multi-point magnetospheric observations to the forefront of space research. The Cluster-II mission elevates the capability of ISEE-1/2 and ISTP missions even further by the ability to determine gradients in three-dimensions as well as differentiate temporal from spatial variations. This capability provides valuable constraints in data interpretation on magnetotail dynamics.

It is evident that the new capability of Cluster can help to advance our understanding of substorm phenomena. Several investigations have utilized Cluster to examine substorm phenomena in the tail [8-12]. In this paper, we explore some substorm features observed by Cluster on 2001 August 22 in some details. This event study is selected based on the availability of global auroral images from IMAGE simultaneous with Cluster observations so that substorm onset and subsequent intensifications can be identified and related to magnetotail disturbances. The main focus of this study are (1) the comparison between dipolarization and plasma flow onsets including an examination of the

related magnetic stress, and (2) the magnetic stress in a plasma flow reversal interval when tailward plasma flow accompanied by southward magnetic field changed to earthward plasma flow accompanied by northward magnetic field. These topics may be used to validate or falsify existing substorm models. In particular, the traditional interpretation of this type of plasma flow reversal is the tailward movement of an X-type neutral line, a key element of the near-Earth neutral line model for substorms. We shall show that there are features in this flow reversal interval that do not fit well with this traditional interpretation. A more plausible interpretation of the flow reversal event discussed here is the shifting role of different activity sites in dominating the plasma behavior at a given point of observation.

## 2. OBSERVATIONS

The ground-station magnetic activity on 2001 August 22 is provided by the AU/AL indices shown in Fig. 1. For the time interval of interest, 0900-1200 UT, the AL index showed the strengthening of auroral electrojet signifying substorm development. The more definitive signature of substorm onset was registered by the IMAGE/FUV imager at  $\sim$ 0920 UT, with a sizeable auroral bulge seen at  $\sim$ 0926 UT, as shown in Fig. 2. Multiple substorm intensifications followed as indicated by the series of auroral images. During this period, Cluster was near its apogee in the magnetotail and crossed the neutral sheet from north to south, as shown in Fig. 3 with measurements from C3. Shown in the panels are the number density, the ion temperature (red trace), the  $x$ - and  $y$ -components of plasma flow, the  $y$ -components of electric field and  $-V \times B$  (cross-product of plasma flow and magnetic field - red trace), the  $x$ - and  $y$ -components of the magnetic field, and the  $z$ -component of the magnetic field. The plasma measurements were taken by CIS/HIA [13], which are used in this entire study. The electric and magnetic field measurements were obtained by EFW [14] and FGM [15], respectively. The measurements show several dipolarizations of the magnetic field and plasma flow reversals, which will be discussed in detail later. It also can be noted that the  $y$ -components of electric field and  $-V \times B$  matched well in general but significant deviations can also be seen intermittently.

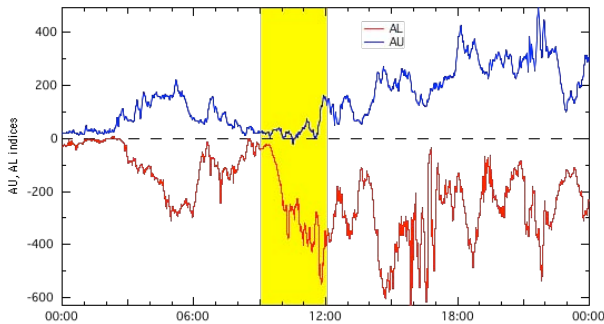


Fig. 1. AU and AL indices for 2001 August 22.

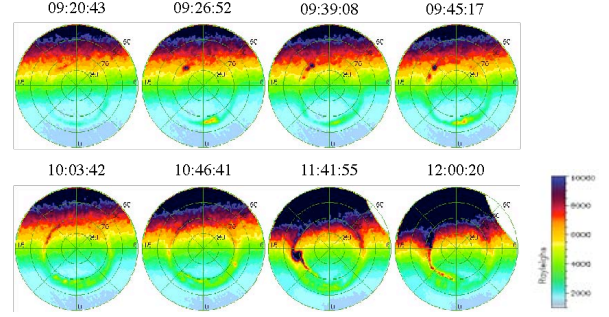


Fig. 2. Global auroral images from IMAGE/FUV.

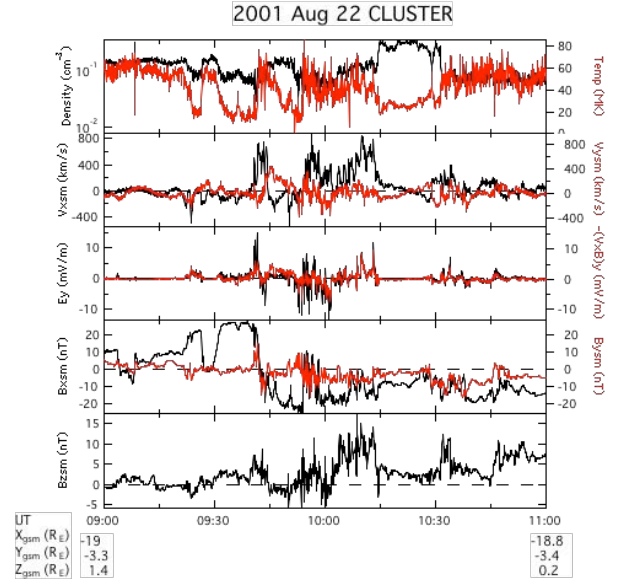


Fig. 3. Overview of magnetotail activity at Cluster.

The dimension of the tetrahedron formed by the four Cluster satellites was  $\sim$ 2000 km, with C3 at its apex south of the tetrahedron base, see Fig. 4. Cluster data for this period have been studied previously with focus on MHD oscillation modes of the plasma sheet rather than substorm features as aimed here [16-18].

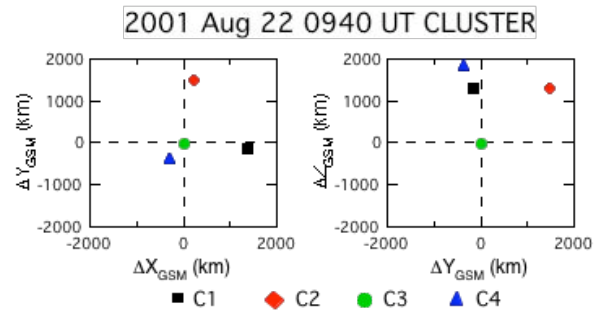


Fig. 4. Relative locations of Cluster spacecraft.

### 2.1 Dipolarization and Plasma Flow

The first set of two dipolarizations is shown in Fig. 5, where measurements from different satellites are given in different colors. Plasma flows shown are components perpendicular to the magnetic field. The geocentric solar

magnetospheric (GSM) coordinates are used throughout this study. Cluster satellites were located at  $\sim(-18.9, -3.3, 1.1) R_E$  and detected dipolarization at  $\sim 0939$  and  $\sim 0940$  UT, which were  $\sim 20$  min after the substorm onset seen by IMAGE/FUV data. Based on the T96 magnetic field model [19], the location of Cluster projected along the magnetic field line to the ionosphere at this time is  $67.5^\circ$  MLat and 1.1 MLT, just poleward of the auroral bulge seen by IMAGE/FUV at  $\sim 0926$  UT. Prior to dipolarization, southward dipping of the magnetic field (i.e., southward  $B_z$  component small compared with the total magnetic field strength) was seen by all four satellites, suggesting plasma sheet thinning before dipolarization onset. Note that dipolarization preceded the occurrence of fast plasma flow. Highly structured electric and magnetic fields, indicative of a turbulent region, were detected after dipolarization. The  $E_y$  component exhibited a kink mode behavior, i.e., duskward in the northern part of the tail and downward in the southern part of the tail. Only the  $E_y$  component is examined in this study because it is the most reliable component from EFW.

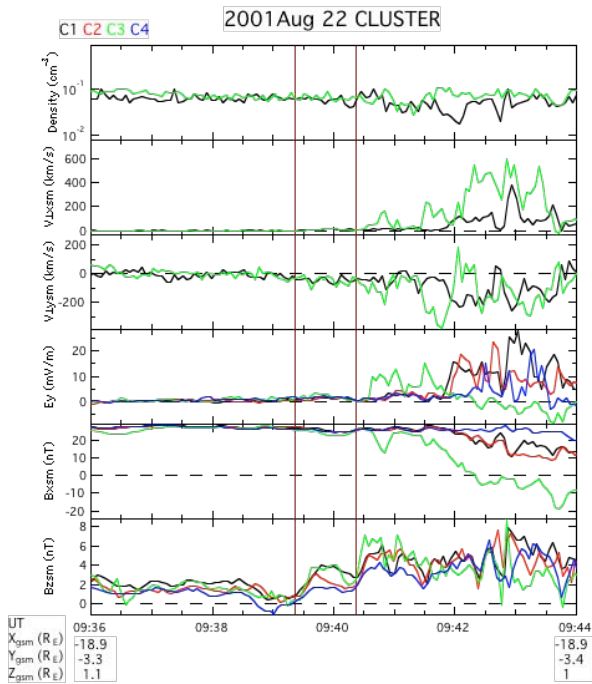


Fig. 5. Dipolarizations at  $\sim 0939$  and  $\sim 0940$  UT.

Another set of two dipolarizations at  $\sim 1002$  and  $\sim 1004$  UT is shown in Fig. 6. The  $E_y$  component was strong prior to dipolarization, maintaining opposite signs in the two halves of the magnetotail as noted in the previous interval in Fig. 5. There was fast earthward plasma flow  $\sim 1$  min before the first dipolarization when the  $B_z$  component was southward. However, at the dipolarization onset times, there was no evidence of fast plasma flows. The dipolarization onsets appeared to be simultaneous at all four satellites with the spin-averaged

data. After dipolarization, the  $E_y$  component diminished its strength.

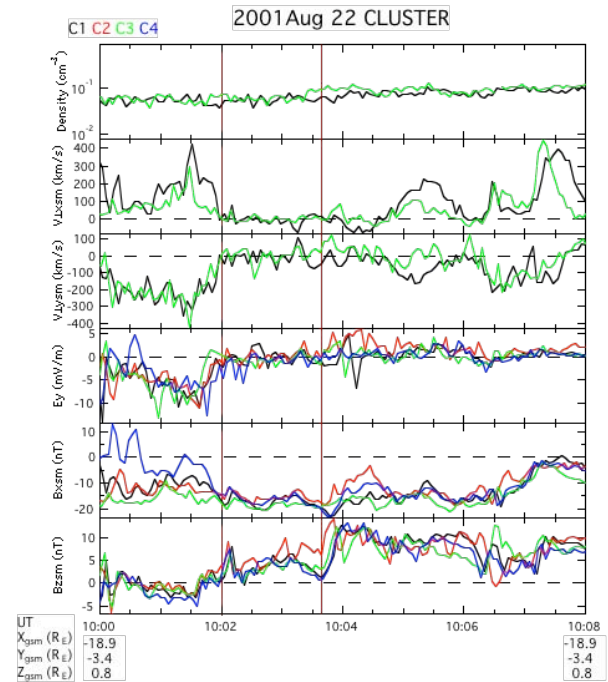


Fig. 6. Dipolarizations at  $\sim 1002$  and  $\sim 1004$  UT.

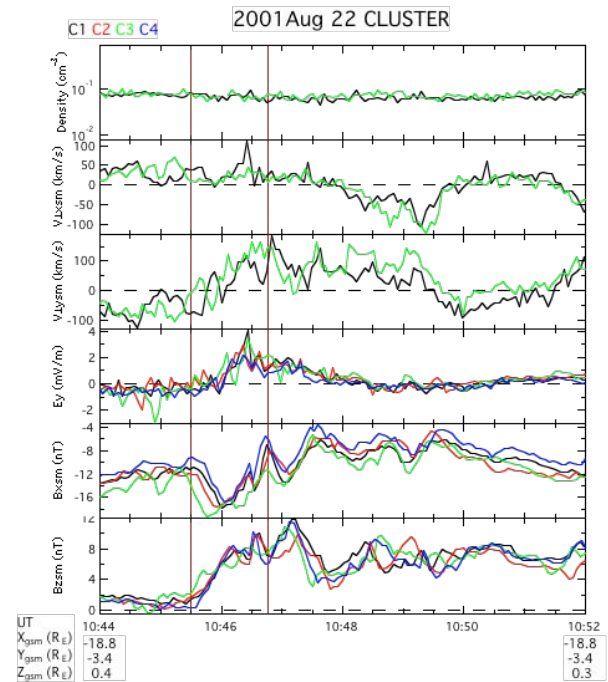


Fig. 7. Dipolarizations at  $\sim 1044$  and  $\sim 1052$  UT.

The final set of two dipolarizations during this interval occurred at  $\sim 1045$  and  $\sim 1047$  UT, as shown in Fig. 7. Around the dipolarization onset times, plasma flow was relatively slow. The  $V_x$  component was below 50 km/s while  $V_y$  was less than 200 km/s. These flow speeds are well below the level qualified for bursty bulk flow, which is often taken to be 300-400 km/s.

In general, it is difficult to use timing analysis to determine the propagation of the dipolarization front since the detailed  $B_z$  profile at each satellite was different from the others. These characteristics suggest that the magnetic field profile changed significantly as dipolarization proceeded from one satellite to the next and/or the front was highly non-uniform spatially. However, for some dipolarizations, the temporal sequence seems to be clear. For instance, for the  $\sim 1004$  UT onset, the sequence was C1-C2-C3-C4, indicating that dipolarization propagated from the earthward-duskward direction. For the  $\sim 1045$  UT onset, the sequence was C2-C1-C3-C4, indicating that dipolarization propagated from the duskward-earthward direction. The observed dipolarization propagation direction is consistent with previous report of tailward propagation of dipolarization [4].

## 2.2 Energetics Associated With Dipolarization

The curlometer technique [20-22] is used to determine the current density and plasma parameters associated with current density in these dipolarizations. Fig. 8 shows the result for the dipolarizations at  $\sim 0939$  and  $\sim 0940$  UT. The panels from top to bottom show, respectively, the three components of current density, the three components of the Lorentz force ( $j \times B$ ), the dissipation/dynamo term  $j_y \langle E_y \rangle$  (the angle brackets indicate averaging over the four spacecraft), and the  $B_z$  component from all four spacecraft as a reference. Prior to dipolarization, the current density was low. This is to be expected since the satellites were not near the neutral sheet region as indicated by the strength of the  $B_x$  component. After dipolarization, the satellites became embedded within the central plasma sheet and observed higher current density than before. The  $(j \times B)_z$  component was most significant and pointed towards the neutral sheet. This implies a strong plasma pressure gradient with the high pressure towards the neutral sheet, which is the expected pressure gradient direction. The  $(j \times B)_x$  component was pointing earthward, consistent with the expected enhancement in the earthward-directed magnetic stress during dipolarization [23]. The dissipation/dynamo term was mainly positive, but it became negative momentarily in the turbulent region after the dipolarization onset.

Fig. 9 shows the same current parameters for the second set of dipolarizations. At the beginning of the interval, the satellites were within the central plasma sheet (the  $B_x$  component was small). The cross-tail current component  $j_y$  was relatively strong. After dipolarization,  $j_y$  became weaker, the three components of  $j \times B$  varied significantly, and the  $j_y \langle E_y \rangle$  term changed sign frequently in the turbulent magnetic field region. Similar behavior of these plasma parameters was seen for the third set of dipolarizations as well.

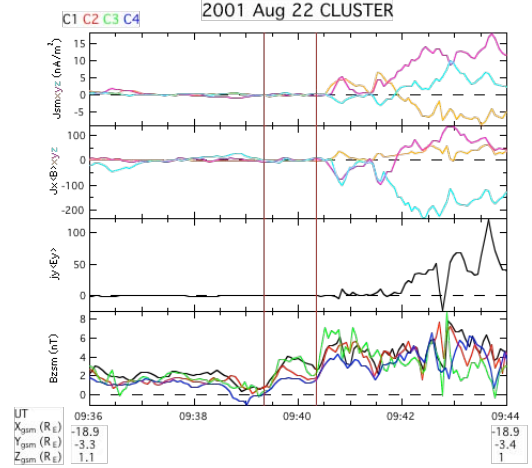


Fig. 8. Current parameters for  $\sim 0939$  and  $\sim 0940$  UT dipolarizations.

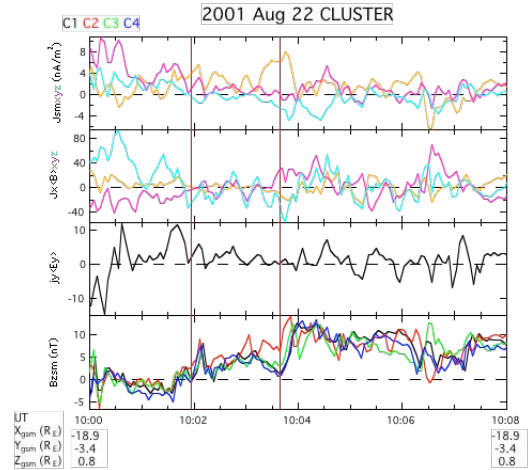


Fig. 9. Current parameters for  $\sim 1002$  and  $\sim 1004$  UT dipolarizations.

## 2.3 Plasma Flow Reversal

There was a plasma flow reversal during this interval of activity, as shown in Fig. 10. The panels show the number density and  $V_x$  component from CIS/HIA on C1 and C3,  $E_y$  and three components of magnetic field from all four spacecraft. The plasma flow reversal at  $\sim 0954$  UT was accompanied by a reversal in the  $B_z$  component. There was a reduction in number density at plasma flow reversal. Note that both C1 and C3 crossed the neutral sheet just around the flow reversal time and yet the number density at the neutral sheet was lower there than earlier in the interval when the spacecraft was further away from the neutral sheet. Thus, there was a substantial loss of plasma in the neutral sheet at that time.

The current parameters during the flow reversal interval are shown in Fig. 11. The  $j_y$  component stayed positive while the  $j_x$  component reversed sign from positive to negative at plasma flow reversal. An interesting feature is the sign reversal of the  $y$ - and  $z$ -components of  $j \times B$  at

flow reversal while the  $x$ -component did not show any reversal. This indicates that the magnetic stress in the  $y$ - and  $z$ -directions changed at flow reversal but not in the  $x$ -direction. The  $j_y \langle E_y \rangle$  term is quite variable, indicating the occasional presence of the dynamo effect in addition to dissipation.

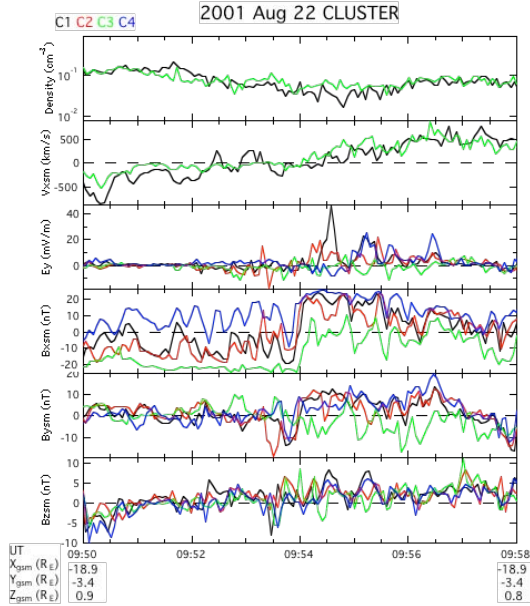


Fig. 10. A plasma flow reversal detected by Cluster.

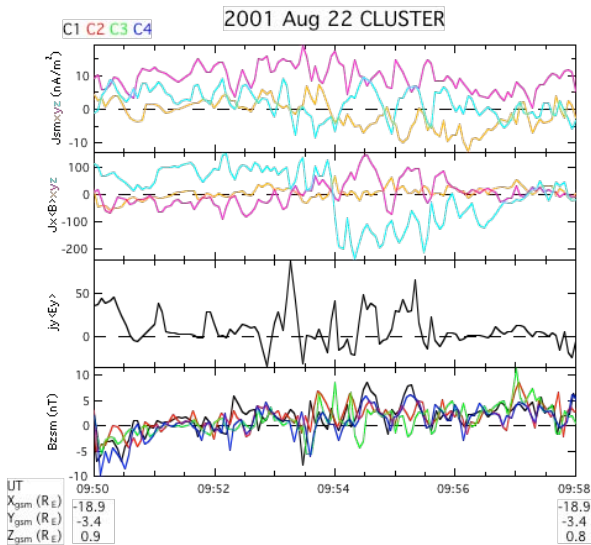


Fig. 11. Plasma parameters associated with current density during a plasma flow reversal.

## 2.4 Breakdown in Frozen-in Condition

During this sequence of substorm activity, there were several brief instances in which the frozen-in condition was violated, i.e.  $E \neq -V \times B$ . Fig. 12 shows three intervals for which the  $y$ -components of  $E$  and  $-V \times B$  are compared together with the three components of the local magnetic field. The first interval encompasses the dipolarization and the data were taken by C3. The

second interval is during the turbulence encountered after dipolarization and the third is the flow reversal period, both showing data from C1. For the first interval, the two quantities matched well prior to dipolarization. However, significant differences occurred afterwards. In the turbulent region (second interval), the two quantities did not agree in general. For the third interval, large discrepancies existed intermittently. Note that the breakdown of frozen-in condition occurred when the magnitudes of  $B$  and  $B_z$  had large values, indicating that these regions are not within the diffusion region around an X-line where the magnetic field is expected to be weak.

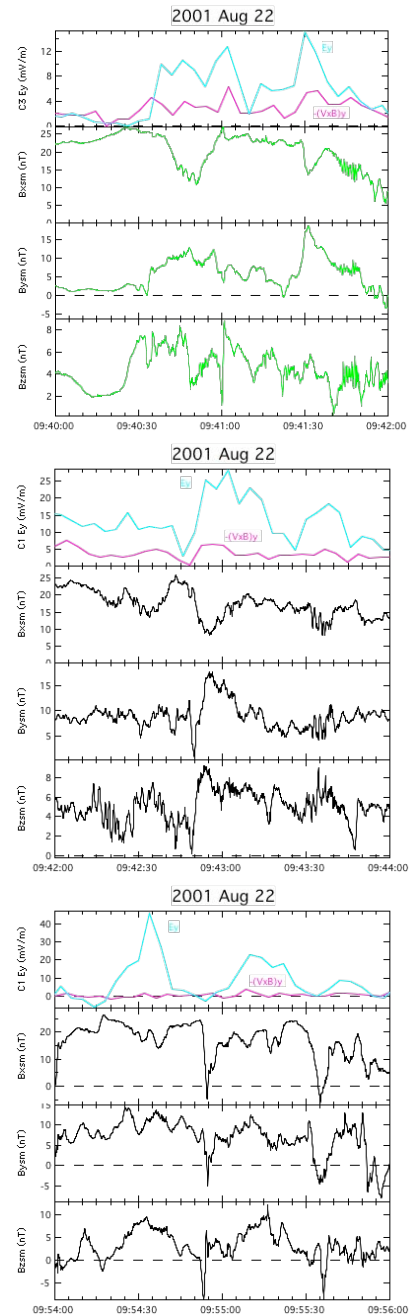


Fig. 12. Three intervals when the frozen-in condition breaks down.

The various terms of the generalized Ohm's law were examined to determine the cause for the violation of the frozen-in condition, i.e., the non-MHD behavior. It is found that the Hall term was the most significant non-negligible term, followed by the electron pressure gradient term. The electron inertial term was negligible. However, the magnitude of the anomalous resistivity term is unclear due to the lack of information on the value of anomalous resistivity. For example, it is found that at 0954:35 UT, in SI units,  $(E+V\times B)_y \sim 4.6\times 10^{-2}$ ,  $(j\times B)_y/Nq \sim 3.8\times 10^{-2}$ ,  $-(\nabla\cdot\mathbf{P}_e)_y/Nq \sim 5.9\times 10^{-4}$ , and  $(m_e/q^2) d(j_y/N)/dt \sim -1.6\times 10^{-7}$ , where the new symbols are the number density  $N$ , the unit electric charge  $q$ , the electron pressure tensor  $\mathbf{P}_e$ , and the electron mass  $m_e$ .

### 3. SUMMARY AND DISCUSSION

The magnetotail behavior during a sequence of substorm activity on 2001 August 22 is examined with data from Cluster satellites that were at  $\sim 19 R_E$  downstream in the midnight sector. Several magnetic field dipolarizations were detected. Dipolarization was episodic, consisting of multiple increases and decreases of the  $B_z$  component. Each dipolarization front was followed by an interval of highly variable electric and magnetic fields as well as variable  $j_y E_y$  sign, suggesting the occurrence of turbulence in the dipolarization process and also the presence of dynamo effect as well as dissipation in the turbulent region. In general, no fast ( $>300$  km/s) plasma flow occurred before dipolarization during this interval. This is evidence for a non-MHD feature, suggesting dipolarization to be related to a kinetic process. A decrease of  $B_z$  prior to dipolarization was also generally seen, suggesting plasma sheet thinning prior to dipolarization. During a plasma flow reversal, there were sign changes in the magnetic stress in the  $y$ - and  $z$ -directions but not in the  $x$ -direction. This feature suggests that the plasma flow reversal arose from shifting dominance of multiple activity sites rather than movement of a single activity site. Note that tailward movement of a near-Earth neutral line should give rise to a changing sign (from negative to positive) of magnetic stress in the  $x$ -direction, which is not seen here. It is unclear how general this result is for flow reversal events. At the very least, however, this finding demonstrates that plasma flow reversal from tailward to earthward is not necessarily due to tailward retreat of a single magnetic reconnection site.

Some of these results summarized above are not new. A decrease of  $B_z$  prior to dipolarization has been reported previously by several researchers [24-28]. Different interpretations were considered, such as plasma sheet thinning from rarefaction wave passage [24], current intensification prior to dipolarization [25, 26], eastward expansion of substorm current wedge [27], and flux rope associated with earthward plasma flow [28]. In the events shown here, decrease of the  $z$ -component of the magnetic field prior to dipolarization was not associated

with plasma flow, thus ruling out the flux rope interpretation. For the interpretation of eastward expansion of the substorm current wedge, it predicts that the  $B_y$  peak will coincide with the zero crossing of the  $B_z$  component. In these events, the associated  $B_y$  component did not show such a trend and thus ruling out this interpretation also. The interpretations of plasma sheet thinning from rarefaction wave passage and current intensification are essentially the same since the former emphasizes the boundary while the latter emphasizes the current density enhancement as a result of plasma sheet thinning. The observed field characteristics (such as southward dipping of the magnetic field) are consistent with these interpretations.

Another quite well known aspect of dipolarization is its episodic behavior [25, 27]. Dipolarizations can occur in succession within 2 min., e.g., events reported in [27]. The lack of plasma flow at the time of dipolarization has also been found previously [29, 27].

The new features reported here are (1) the variability in the dynamo/dissipation effect within the turbulent region following the dipolarization, (2) the absence of the sign reversal in the  $x$ -component of the Lorentz force in association with plasma flow reversal - other components reversing in sign instead, and (3) the breakdown of the frozen-in condition in several regions associated with the dipolarization activity. These regions do not exhibit low magnetic field strength expected for the diffusion region around a magnetic reconnection site. Furthermore, these regions had large  $B_z$  component, suggesting that they are not associated with the separatrix layer of magnetic reconnection, which should coincide with the plasma sheet boundary, where the  $B_z$  component would be small. The violation of the frozen-in condition may be due to a kinetic plasma process, such as the kinetic ballooning instability [30, 31], the current-driven electromagnetic ion cyclotron instability [32], the cross-field current instability [33], or the entropy antidiffusion instability [34], that are related to current disruption and dipolarization. Large amplitude high frequency waves expected from the excitation of the cross-field current instability have been observed by Geotail in high-time resolution wave data during dipolarization without accompanying plasma flow [27]. This aspect will be considered in future extension of this work.

Many articles in the literature on magnetotail dynamics interpreted observations in terms of magnetic reconnection only without considering alternative explanations. However, many of the observed features are basically products of ions decoupling from electrons in the plasma, a natural outcome of a kinetic plasma process. Therefore, the observed features are not unique to magnetic reconnection alone. For example, the generation of the so-called Hall current system (together with the quadruple out-of-the-plane magnetic

perturbations) associated with magnetic reconnection [35] is essentially (and more appropriately called) a meridional current system when ions become unmagnetized while electrons are not. It is pointed out [36, 37] that such a current system can be generated by the onset of a kinetic plasma process like the cross-field current instability. A bipolar electric field signature, as seen in some dynamic magnetotail events [11], can also be created as a result. It is prudent for future investigations of magnetotail dynamics to distinguish these models by identifying observational features that are not common among these different models.

#### 4. ACKNOWLEDGMENT

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